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Automated extraction of single H atoms with STM: tip state dependency

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Abstract
The atomistic structure of the tip apex plays a crucial role in performing reliable atomic-scale surface and adsorbate manipulation using scanning probe techniques. We have developed an automated extraction routine for controlled removal of single hydrogen atoms from the H:Si (100) surface. The set of atomic extraction protocols detect a variety of desorption events during scanning tunneling microscope (STM)-induced modification of the hydrogen-passivated surface. The influence of the tip state on the probability for hydrogen removal was examined by comparing the desorption efficiency for various classifications of STM topographs (rows, dimers, atoms, etc). We find that dimer-row-resolving tip apices extract hydrogen atoms most readily and reliably (and with least spurious desorption), while tip states which provide atomic resolution counter-intuitively have a lower probability for single H atom removal.

Keywords: scanning probe microscopy, STM, atomic manipulation, hydrogen desorption, tip state

(Some figures may appear in colour only in the online journal)

1. Introduction

In scanning probe microscopy (SPM), the ability to image and perform atomic and molecular manipulation has unsurprisingly been shown to be critically dependent on the state of the tip. Since their invention in the eighties, these instruments have been invaluable for investigating and manipulating surfaces and adsorbates at the atomic scale. An especially important class of atomic manipulation tool is available via the scanning tunnelling microscope (STM), which has the capability to perform hydrogen resist lithography on H:Si (100) surfaces [1–4] with single atom specificity [5]. This technique has been used to fabricate not only prototype but functional electronic devices, including quantum dots [6, 7], conducting wires [8], and single atom transistors [9], and has been elegantly exploited to generate and manipulate artificial molecular orbitals (arising from dangling bond (DB) interactions) at the Si(100) surface [10]. One of the main advantages of STM-induced depassivation of this type is that it allows atomic scale imaging in concert with the patterning process, thus offering a high degree of operational control.

With the ultimate goal of autonomously fabricating atomic scale structures, we have implemented the following protocols. In order to automate and control the H:Si(100) patterning process with single-atom specificity, a feedback-controlled lithography (FCL) approach based on the technique developed by Lyding et al [5] was implemented. By actively monitoring the STM feedback signal and controlling the desorption parameters during patterning, individual atoms can selectively and ‘autonomously’ be desorbed. Similar protocols for automated atom-precision lithography have been developed by Randall and co-workers at Zyvex Labs...
the resulting control system, ZyVector, is now commercially available. We show that for the variant of the FCL technique we have adopted, robust detection of single desorption events is possible. More importantly, the automation protocols we have developed enable the systematic generation and analysis of thousands of hydrogen desorption events, which in turn have provided key insights into the role that the atomistic structure of the tip apex plays in the extraction process.

The STM desorbs hydrogen atoms through inelastic electron scattering mechanisms [2, 3, 12, 13]. At conditions of relatively high current and low electron energy tunneling, such as those used in our experiments, the desorption mechanism involves multi-quantum vibrational excitations of the Si–H bond [2, 12, 14–16]. At higher electron energies electrons possess sufficient energy to directly promote an electron from the Si–H bonding state to its corresponding antibonding orbital. Previous studies have demonstrated that the basic desorption process involves only a single Si–H bond which is vibrationally excited by tunnelling electrons [2, 17]. Although the H:Si(100) surface [18–21], and the desorption mechanism [2, 3, 12, 13], have been studied extensively, the influence of the tip state on the desorption process has hitherto not been examined in any depth. For example, while the desorption efficiency has been analyzed for different bias and current parameters [22], these experiments did not take the state of the tip into account. Theoretical studies of atomic force microscope-based atomic manipulation [23–25] clearly, and unsurprisingly, show that the atomistic detail of the tip termination plays a crucial role in the process. (Unlike STM-tip-induced desorption, however, AFM-based manipulation on Si(100) surfaces is mediated by direct chemical interaction.)

In this work, we examine the influence of the atomistic structure of the tip apex on the desorption efficiency for STM-induced H atom extraction. Through a detailed, systematic analysis of thousands of H atom desorption events, we show that tips that produce row-resolved STM topographs yield single atom desorption events more reliably than tip apices which generate either dimer-resolved or atomically-resolved images. Moreover, row-resolving tip apices also generate the smallest number of spurious desorption events. (This dependence of spurious hydrogen desorption on the tip state was postulated previously by Ballard et al [26]. Here we provide strong evidence to support their hypothesis.) This sensitivity of H atom extraction to the precise atomistic structure of the tip is a critical factor in ensuring the reproducibility and reliability of atomic scale engineering of hydrogen-passivated semiconductor surfaces.

2. Methods

Silicon samples cut from Si(100) wafers (n-type, 6 mΩ cm) were used throughout our experiments. Samples were prepared in an ultra-high vacuum (UHV) chamber with a base pressure of $5 \times 10^{-11}$ mbar or less. Low base pressures were obtained by baking the UHV chamber at 145 °C twice over prolonged periods (three to four days). The samples were outgassed for a minimum of 12 h at 600 °C. Clean Si(100) 2 × 1 samples were obtained by flash heating to 1150 °C for 10 s, followed by gradual cooling from 900 °C (at a rate of $\sim 1 {$°}C$ s$^{-1}$) to room temperature. For hydrogen passivation, the sample was exposed to atomic hydrogen, which is obtained by cracking molecular hydrogen using a Specs GmbH thermal gas cracker, positioned approximately 10 cm from the sample. H₂ was introduced into the chamber, through the gas cracker, until a pressure of $1 \times 10^{-6}$ mbar was reached. The sample was exposed for 1 min (i.e. 60 Langmuirs). During exposure, the sample was held at a temperature of 360 ± 20°C in order to achieve hydrogen-terminated Si(100)-2 × 1 surfaces and avoid other reconstructions (such as $(3 \times 1)$).

Measurements were performed on an Omicron variable-temperature (VT) STM-AFM system with a MATRIX control system. All experiments were performed at room temperature (RT). Tungsten tips were used for both STM imaging and hydrogen extraction. These were electrochemically etched and subsequently annealed to ≤1000 °C in the UHV chamber, in order to remove the oxide layer, before STM use. In order to randomize our statistics, the tip apex was routinely prepared during STM imaging using voltage pulsing and controlled tip crashing. Atom tracking [27] was used to eliminate drift during all automated extraction routine trials, unless otherwise stated. All STM images presented in this article were obtained with a sample bias voltage of 1.6 V and a constant current of 10 pA.

3. Results

Figure 1 shows a flow diagram of the automated DB generation routine. Upon activation of the automation routine, the tip was positioned at the desired location for atom extraction, and the initial parameters were subsequently changed to the desired desorption parameters $V_D$, $I_D$. Upon reaching the desorption parameters, the $Z(t)$ feedback signal is then used to monitor in real-time for a jump in $Z$. A threshold jump height of 0.3 Å was found most suitable for detecting a desorption event, as higher thresholds potentially lead to desorption events not being detected, while smaller thresholds increase the probability of false detections. If a step in $Z$ ($Z_{jump} \geq 0.3$ Å) is detected, the desorption parameters subsequently rebound to the initial scan parameters. For step detection, a 10 ms integration time was found to be most appropriate as larger integration times could lead to multiple desorption events. Figure 2 shows a typical $z$-jump detection event, including the full voltage- and current-ramp-related changes in the $z$ position, for generation of a single DB.

The automated extraction routine can be used for arbitrary patterning with high fidelity as shown in figure 3, where the extraction routine is used to atomically generate a $4 \times 4$ array of DBs. The reliability of the automated extraction routine was evaluated by comparing step detections from the routine with the true outcome as observed from STM imaging. The results are plotted in figure 4, for 1126 trials. For
each trial, the extraction routine is run for 60 s using similar desorption parameters. Desorption events are classified based on whether any desorption occurs directly underneath the tip (primary) or if none occurs directly underneath the tip. Desorption events with either an uncertain number, or more than four DBs generated, are collectively called clusters. We find that our step detection algorithm reliably detects any type of desorption, whether they occur directly underneath the tip (primary) or not, but elsewhere (secondary). Single atom extraction events are detected with a 90% accuracy. Other unwanted events, such as multiple desorption events, or secondary desorptions, are similarly reliably detected. Unfortunately, the automated desorption algorithm does, unavoidably, occasionally register false jumps. This is mainly due to a combination of noise in $Z$ and a low $Z$-threshold value. As discussed below, the tip state plays a key role in triggering false detection events.

It is well known that the structure of the tip can change during a desorption event as a result of tip–surface interactions \[4\]. To investigate the desorption efficiency under practical conditions, data containing tip changes during desorption were not excluded. Only a small subset (8.9%) of the total number of trials involved tip-state changes. All tip imaging states were classified based on the image observed prior to desorption. Only data that could not be classified into the four tip imaging states (row, dimer, atomic, asymmetric) prior to desorption were excluded (constituting 1.7% of trials).

The dependence of the desorption efficiency on the atomic structure of the tip apex was determined by grouping the tip states into a set of broad classes. Figures 5(a)–(d) show a selection of observed tip states, and their respective grouping into the following classes: (a) rows, (b) dimers, (c) atomic, and (d) asymmetric.

The desorption efficiency is plotted in figure 6 as a function of voltage, for each respective type of observed tip states [(a) row, (b) dimer, (c) atomic, and (d) asymmetric]. All tip states that could be classified into these groupings (2023 out of 2074 trials) were included in our analysis. Each data value in the figure corresponds to an experimental run of the automated extraction routine. In order to gauge the desorption efficiency as a function of voltage, all data were acquired by keeping the desorption parameters fixed (1.2 nA, $V_{fix}$) for a given bias for the same amount of time (60 s). All

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**Figure 1.** Flow diagram showing implementation of automated DB generation. The tip is initially positioned using predetermined desorption parameters and coordinates, followed by a change from scanning conditions to desorption parameters ($V_D$, $I_D$, $t$). The feedback signal, $Z(t)$, was then monitored in real time. Following detection of a jump in $Z$, or an elapsed time $t$, the parameters were subsequently restored to the imaging conditions. The desorption procedure was then repeated for all input coordinates.
experimental trials performed at each voltage value \( (V_{\text{fixed}}) \) were therefore normalized to 1. Data where false detections occurred before 60 s were excluded.

We observe that the onset of desorption follows the general tip classification rule: \( V_{\text{row}} \geq V_{\text{dimer}}, V_{\text{asymmetric}} \geq V_{\text{atomic}} \) for a given bias, respectively. This translates into a desorption probability that follow the same rule: \( P_{\text{row}} \geq P_{\text{dimer}}, P_{\text{asymmetric}} \geq P_{\text{atomic}} \). These results are further corroborated by similar findings using other STM tips and using different FCL automation routines (not included).

4. Discussion

Our findings show that row-resolved tip states produce a single desorption event more reliably than other tip states, and
that other unwanted reactions (such as producing clusters) are less likely for the row-resolving tip apex than for other states. The reason for this remains as yet unclear but an obvious candidate explanation relates to the density of states of the tip apex. Although the STM image contrast mechanism has been investigated for bare [28, 29], and passivated Si(100) surfaces [30, 31], this has been in the context of the density of states and band structure of the substrate; the influence of the atomistic structure of the tip on the contrast mechanism has not yet been investigated.

For our desorption conditions (electron energies < 4 eV), the dominant desorption mechanism is attributed to inelastic scattering of the tunneling electrons with the Si–H bond [2, 12]. Using this model, the tunneling probability is dependent on several factors, including, of course, the local density of states (LDOS) of the tip. It is therefore reasonable to speculate that for the H:Si(100) surface, tip states producing row-like STM topographs might have a greater density of states (in the energy window defined by the tip–sample bias) than other classes of tip apex. This will directly affect the available flux of electrons and, thus, the desorption probability.

Although early theoretical studies suggested that while the electric field under the tip is insufficient to directly induce H atom desorption, it could nonetheless play a significant role via spatial localization of the excitation [14] due to a Stark effect. However, subsequent experiments, where neighboring hydrogen atoms were extracted, show that interdimer extraction is more favorable than intradimer extraction, suggesting that the presence of an electric field does not play a significant role in the desorption process [17]. Our results support these latter findings. Moreover, we find that the distribution of different types of neighboring desorption events is also tip-state dependent, where intradimer desorption occurs more frequently for atomically resolved STM topographs. We are planning a systematic combined density functional theory (DFT)-non-equilibrium Greens function study in order to elucidate the origin of the higher desorption efficiency of the row-resolving class of tip apices (and to ascertain the extent to which the H desorption process is mediated by the structure and density of states of the tip apex).

Our findings have shown that extracting single H atoms is best achieved for row-resolving tip imaging states and low bias voltages. Additionally, at these conditions, undesirable side effects such as multiple or secondary desorptions are also minimized. The low bias voltage probably plays a significant role in limiting the number of unwanted desorption events as the yield is highly dependent on the bias voltage and electron dosage [2–4]. Based on our findings we conclude that the best strategies for controlling the automated patterning process can be obtained by (I) carefully selecting the proper tip geometries for desorption, (II) controlling the desorption yield (affected by desorption parameters, tip–sample distance), and (III) through reliable detection of H extraction by either minimizing/excluding noise influences or improving the step detection algorithm.

Figure 5. Classification of tip imaging state into: (a) rows, (b) dimers, (c) atoms, and (d) asymmetrical. The STM images have been rotated to align the direction of dimer rows. Line profiles (e) along and (f) perpendicular to rows from positions indicated in (a)–(d). Note that some of the profiles have been shifted by an offset of 20 pm for clarity. Scale bars correspond to 0.77 nm (i.e. the dimer row separation on Si(100)).
5. Conclusions

We have implemented an automated extraction routine capable of implementing and detecting single H atom desorption from H:Si(100) surfaces with high reliability. The routine was used to examine the influence of the tip state on the desorption efficiency by desorbing atoms using fixed desorption parameters, for a fixed time. Tip states were classified into four broad types: row-resolving, atom-resolving, dimer-resolving, and ‘asymmetric’ and their desorption efficiencies compared. We find that tips that produced row-resolved STM topographs desorb H atoms most efficiently. Our results are used to propose strategies for controlling the accuracy of single H atom extraction. Future work will focus on elucidating the origin of the tip state dependency of the desorption process via a combined density functional theory (DFT)-Green’s function approach.

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