ABSTRACT

This article aims to offer a critical overview of selected literature on capacitive and non-capacitive faradaic charge storage. It is particularly relevant to the concept of pseudocapacitance that is generally described as a result of fast surface faradaic processes. In general, faradaic processes represent electron transfer reactions at the interface between an electrode and its contacting solid or liquid electrolyte phase that is able to accept or donate electrons. Obviously, not all faradaic processes can be associated with pseudocapacitance. The question is how to differentiate pseudocapacitance related faradaic charge storage from the others. Therefore, attempts have been made to apply the band model for semiconductors to account qualitatively for the origin of pseudocapacitance. Capacitive and non-capacitive faradaic processes are then proposed to define and differentiate different charge storage mechanisms in supercapacitor and battery. On the other hand, the unequal electrode capacitance approach and the use of Ca\(^{2+}\) in aqueous electrolytes are discussed in relation with enhanced energy capacity of supercapacitors. In addition, the principle of supercapattery as a hybrid device is explained with recent literature examples.

Key words: Capacitive charge storage; Faradaic process; Supercapacitor; Supercapattery; Unequal electrode capacitance.
1. Introduction

Since the beginning of human civilisation, the use of energy has been an important
means of survival. From burning wood to nuclear energy, the search for and the utilisation of
various energy forms form an important part of human history. In recent years, due to the fast
depletion of fossil fuel reserves and the increasing negative environmental impact of CO₂
emission, political and economic emphases have been placed on the rapid development of
alternative energy technologies, particularly through solar, wind and wave generation.
However, the energy derived from these renewable sources must be stored and supplied
efficiently in order to be considered a viable competitor to traditional non-renewable options
[1]. Energy storage devices including supercapacitor, battery and supercapattery are such
devices that are able to store charges in a fast and efficient way and hence help harvest and
convert renewable energy to a usable form.

Electrical energy can be electrochemically stored in two fundamental ways: (1) in
solid electrode materials relying on fast charge separation and/or chemical reactions of the
materials and (2) in liquid electrolytes flowing through a battery of electrochemical cells. The
latter is typically represented by redox flow battery. It is unique because energy storage is
achieved by reduction and oxidation of redox couples in different electrolytes separated by a
membrane and circulated in their own respective spaces [2]. However, due to the requirement
for storage and transport of a large volume of the electrolytes, the application for redox flow
batteries is limited to stationary needs at present. Great efforts have therefore been made to
deply and comprehensively understand charge storage processes in solid electrode materials
in order to develop versatile energy storage devices (e.g. supercapacitors) and improve their
performance, which is also of interest of this study.

In general, electrical charge can be stored in the surface and/or solid phases of the
electrodes relying on the following storage mechanisms: double layer charging, faradaic
processes, or a combination of both. For instance, the charging-discharging process in an activated carbon electrode is due to charge adsorption or accumulation at the interface between the electrode and the electrolyte, giving rise to the electric double layer (EDL) capacitance. This process is largely electrostatic and non-faradic in nature. It means that ideally, no electron transfer takes place across the electrode/electrolyte interface and the storage of electric charge and energy involves no chemical and physical changes inside the solid phase of the electrode [3]. For the charge storage in a faradaic process, it is generally accepted that large pseudocapacitance arises from electrosorption of ions accompanied by surface redox reactions in which electrons transfer occurs crossing the interface of the current collector and active material.

It is however necessary to point out both EDL and pseudocapactive storage processes are capacitive in nature. In experimental terms, they both offer rectangular cyclic voltammograms (CVs) and linear (or triangular) galvanostatic charge and discharge plots (GCDs) [4]. On the other hand, when the redox active material on electrodes undergoes a reversible or quasi-reversible electron transfer reaction at well separated or isolated sites, current peaks forms on the CV and non-linear potential (or voltage) variation occurs on the respective GCDs. Such electrochemical features are the same or comparable to those of recharge batteries, and should not be used for a capacitance measurement. Thus this type of charge storage is recognised as a non-capacitive faradaic process here.

Consequently, one should not associate all these well-known faradaic processes with pseudocapacitance. The question is how to differentiate pseudocapacitance related faradaic charge storage from the others, namely, the non-faradaic charge accumulation in the EDL and the non-capacitive battery-like behaviours of some electrode materials. Herein, the authors propose to define and differentiate the aforementioned charge storage mechanisms according to Fig. 1. Attempts have also been made to account for the origin of pseudocapacitance using
the band model for semiconductors. Further, the recent developments in gaining improved supercapacitor performance are introduced and discussed in terms of modifications of aqueous electrolytes and an unequal electrode capacitance approach in laboratory.

![Fig. 1. Schematic correlation between EDL capacitor, pseudocapacitor, battery and supercapattery (= hybrid of supercapacitor and battery) in terms of capacitive and faradaic charge storage processes.](image)

In order to understand how supercapacitors store large amount of energy and discharge it for high power applications, it may be easier to start with a review on the fundamentals of the conventional capacitors that are widely used in electronic devices. In principle, the amount of charge \(Q\) stored is in proportion to the voltage \(U\) applied across the capacitor. The proportionality is called capacitance \(C\) which links \(Q\) and \(U\) according to Eq. (1):

\[
Q = CU \tag{1}
\]

\(C\) is a constant determined by the dielectric constant of the dielectric medium and proportional to the ratio of the area of the electrode/dielectric interface and the separation distance between the two electrode plates of the capacitor. When the applied voltage varies linearly with time \((t)\), i.e. \(U = U^0 + vt\), where \(U^0\) is the starting cell voltage which may be zero and \(v\) is the linear variation rate of voltage, Eq. (1) can be mathematically converted into Eq. (2), which can be further simplified to Eq. (3) by correlating the current \(I\) with \(v\) [4].
\[ \frac{dQ}{dt} = C \cdot \frac{dU}{dt} + U \cdot \frac{dC}{dt} = C \cdot \frac{dU}{dt} \quad (2) \]

\[ C = \frac{dQ / dt}{dU / dt} = \frac{I}{v} \quad (3) \]

Therefore, the current flowing through a capacitor is in a linear relationship with \( v \), but independent of \( U \). For a constant \( C \), Eq. (3) gives the rectangular \( I \) – \( U \) plots as shown in Fig. 2a, which is also called cyclic voltammograms or CVs. In Fig. 2a, the current, \( I \), not only increases proportionally with the increase of \( v \), but also follows the direction of voltage variation, i.e. \( v \) is positive when voltage increases and it is negative when voltage decreases. The CV is very useful for evaluation of capacitive behaviour of a device or electrode made from a synthetic pure or composite material.

In practice, galvanostatic charging-discharging (GCD) is also used to calculate the capacitance of a device or electrode material. Under a constant current, according to Eq. (3), the charging process of the capacitor reflects a constant rate of voltage increase and vice versa, the discharging relates to a voltage decrease in a constant rate. As a result, a triangular curve is observed as shown in Fig. 2b where the voltage of the capacitor is plotted against the:

**Fig. 2.** (a) CVs at indicated linear variation rates of voltage and (b) GCD plots at indicated constant currents as derived from Eq. (3) for a 50 mF capacitor with \( U_{\text{max}} = 5 \) V. In the GCD plots in (b), \( t_{\text{max}} = U_{\text{max}} C / I_0 \) [4].
time during a cycle of constant current charging and discharging.

When the capacitor is charged, a voltage, $U$, will build up across the two electrodes. The amount of energy ($\varepsilon$) stored and the power output ($P$) from the capacitor can be determined through a series of calculations below. At first, the work, $d\varepsilon$, done in a short time to move a small quantity of charge, $dQ$, to be accumulated at the electrode/dielectric medium interface can be calculated from Eq. (4) by considering $d\varepsilon = UdQ$. If the heat loss during charging is insignificant, the total energy stored in the capacitor is then determined after integration. It is worth noting that a practical capacitor has always a maximum tolerable voltage, $U_{\text{max}}$, beyond which the dielectric or ionic medium will break down (or decompose). Thus there is a maximum energy capacity, $\varepsilon_{\text{max}}$, that can be correlated to $U_{\text{max}}$ and the capacitance, $C$, which is the property of the capacitor (or electrode) material used.

\[ d\varepsilon = U \cdot dQ = \frac{Q}{C} \cdot dQ \]  
\[ \varepsilon_{\text{max}} = \int_0^Q \frac{dQ}{C} = \frac{1}{2} \cdot \frac{Q^2}{C} = \frac{Q \cdot U_{\text{max}}}{2} = \frac{C \cdot U_{\text{max}}^2}{2} \]  

Secondly, it is to determine the power output, $P$, using the time needed to fully discharge the capacitor as expressed in Eq. (6):

\[ P = \frac{\varepsilon}{t} = \frac{C \cdot U^2}{2 \cdot t} \]  

In fact, Eq. (6) can be further developed into Eq. (7) by taking the equivalent series resistance (ESR) and the working load ($R_L$) into account since the ESR is the unavoidable nature of an electric power source. The power transferred from the source to the load can also be derived from $P = iU = i^2R_L$.

\[ P_{\text{max}} = \left( \frac{U_{\text{max}}}{R_L + \text{ESR}} \right)^2 R_L = \frac{R_L \cdot U_{\text{max}}^2}{(R_L + \text{ESR})^2} \]  

where $U = I \cdot (R_L + \text{ESR})$ with the current passing through the circuit shown in Fig. 3. It is
obvious that the maximum power, $P_{\text{max}}$, can be reached at $R_L = \text{ESR}$, leading to Eq. (8) which shows $P_{\text{max}}$ is a function of $U$ and ESR, but independent of $C$, although $C$ determines the amount of energy stored in the capacitor. Furthermore, by bringing Eq. (8) into Eq. (6), the shortest discharging time, $t_{\text{min}}$, can be obtained, which is practically useful and important for device design, such as supercapacitor supported systems. However, Eq. (9) has currently not gained enough recognition or attention in terms of reporting data and analyses. This is somehow unfortunate because highlighting the maximum power of supercapacitors against that of batteries without mentioning the lasting time the device can work at the high power may confuse or even mislead readers and customers who are not well informed.

$$P_{\text{max}} = \frac{\text{ESR} \cdot U^2}{(\text{ESR} + \text{ESR})^2} = \frac{U^2}{4 \cdot \text{ESR}} \tag{8}$$

$$t_{\text{min}} = \frac{C \cdot U^2}{2 \cdot P_{\text{max}}} = \frac{C \cdot U^2}{2 U^2/(4 \cdot \text{ESR})} = 2 \cdot C \cdot \text{ESR} \tag{9}$$

**Fig. 3.** The electric circuit for derivation of $P_{\text{max}}$ and $t_{\text{min}}$ (cf. Eq. (7) - (9)). $R_L$ is the resistance of the working load, $U$ and ESR are the voltage and the equivalent series resistance of the power source, respectively.

2. Non-faradaic capacitive storage

The capacitance of a conventional capacitor typically ranges between $10^{-6} - 10^{-2}$ F, therefore the energy stored in the capacitor is too small for meaningful practical uses. For
example, for a 50 mF capacitor with an applied voltage of 100 V, the energy stored is only 250 J. Hence, in recent years, supercapacitors, also known scientifically as electrochemical capacitors, become an important development in the field of energy storage and conversion. The capacitance of commercial supercapacitors can commonly reach to $10^2$ - $10^4$ F. However, this huge difference in capacitance does not change the fact that a supercapacitor is still a capacitor. There is no doubt that the behaviour of a supercapacitor should also be governed by all the equations discussed in the previous section and be the same as those in Fig. 2, showing the rectangular CVs and triangular GCD plots.

2.1. Supercapacitors with double layer capacitance

EDL capacitors are the first generation supercapacitors in which charge storage occurs in the electric double layer at the electrode/electrolyte interface and relies mainly on charge separation. Although this storage mechanism is the same as that in the traditional electrolytic capacitors, the specific capacitance ($C_s, \text{ F g}^{-1}$) increases drastically by taking advantages of the high porosity and large specific surface area ($\text{m}^2 \text{ g}^{-1}$) of the active electrode materials, e.g. activated carbon. Fig. 4 illustrates the displaced ions in the electrolyte in the interconnected pores of the electrode and the balancing charges (i.e. excess of electrons or holes) on the wall surfaces of the pores that are formed between the packed activated carbon particles. Assuming that the electrode/electrolyte interface has true capacitance of $0.1$ – $0.3$ F m$^{-2}$ and the specific surface area of activated carbons is $1000$ m$^2$ g$^{-1}$, the specific capacitance, $C_s$, of activated carbon can be calculated to range from 100 to 300 F g$^{-1}$. However, in many cases, the specific capacitance of activated carbons is usually about 100 F g$^{-1}$ [5]. The cause is largely because the high porosity and specific surface area make activated carbons weaker and less conducting. The other issue is that not all the internal surface area, such as those of the wall of micro-pores, can be accessed by ions in the activated carbon for charge storage.
Fig. 4. Schematic illustration of charge storage mechanism in an EDL capacitor using porous and high surface area activated carbon particles as the electrode.

Therefore, nanomaterials, such as carbon nanotubes (CNTs) and graphenes, have been suggested as a good solution to solve the problems [6-8]. In addition to the extra high specific surface area and the construction of three dimensional structures, CNTs and graphene can also benefit the electron conduction through carbon structures and the ion conduction through the liquid electrolyte contained in the carbon structures [9].

In particular, CNTs have been widely studied and recognised as suitable materials for applications in supercapacitors with the acknowledgement of the novel structure, narrow pore size distribution, low resistivity and high mechanical strengths [10-13]. Comparing with the porous activated carbon in the upper bound, CNTs have exhibited higher capacitance, e.g. 102 F g⁻¹ for multi-walled CNTs [14], and 180 F g⁻¹ for single-walled CNTs [15]. It has also been reported that a single-walled CNT based supercapacitor showed a promising power capability of 20 kW kg⁻¹ with a maximum energy capacity of around 7 Wh kg⁻¹ at 0.9 V in a solution of 7.5 N KOH [15]. On the other hand, factors that affect the performance of CNT
based supercapacitors have been extensively researched from the aspects of structural
properties (e.g. diameter, length and pore size), graphitisation by heating and modification via
chemical activation, functionalization and surface treatment [6]. Among these, it was found
that boiling CNTs in a mixture of concentrated acids was one of the most effective methods
to modify CNT surfaces and introduce functional groups which play an important role on the
capacitance improvement [16, 17]. However, effects of the electrolytes on the capacitance of
supercapacitors have attracted relatively less attention [18-20]. It is acknowledged that
cations, anions and pH of the electrolyte may all affect the performance of a particular
electrode material. In this paper, CNTs and activated carbon electrodes are investigated in
correlation with the concentration and nature of ions, temperature, and the CNT surface
functional groups. The results are particularly interesting, confirming different affinities of
cations in the electrolyte (i.e. Ca\(^{2+}\) and K\(^{+}\)) towards the oxy-groups on CNT surfaces and
workability of this type of supercapacitors at temperatures below −60 °C.

2.2. Effects of electrolyte

In the experiment, the effect of the electrolyte concentration on the electrochemical
behaviour was investigated by performing cyclic voltammetry in various KCl aqueous
solutions using a 5 mm dia. graphite disc electrode coated with a thin layer of activated
carbon (Haycarb, loading: 0.2 mg or 1.0 mg cm\(^{-2}\)) as the working electrode. A graphite rod
was used as the counter electrode and the Ag/AgCl couple used as the reference electrode.
The resulting CVs are shown in Fig. 5. These results agreed to certain extend with some
reported observations in literature that EDL capacitance increased with the increase of the
electrolyte concentration [21, 22], i.e. the CV current kept increasing until the concentration
of the KCl solution reached to 1.0 M (see Fig. 5a). Interestingly, the CVs observed with
further increasing the concentration to 3.0 M had shown insignificant changes in current as
Fig. 5. (a,b) CVs of the Haycarb carbon electrode (loading: 1 mg cm\(^{-2}\)) in KCl aqueous solutions with different concentrations as indicated. Potential scan rate: 100 mV s\(^{-1}\).

shown in Fig. 5b. It was likely that the active sites on the surface of Haycarb carbons were saturated by the ions in the electrolyte while recording the previous CVs and became inaccessible even though the ion concentration was significantly increased.

The capacitance improvement from the surface modification with functional groups is well demonstrated by the significant current increase in the CVs recorded in 1.0 M KCl using a 6.5 mm dia. graphite disc electrode coated with acid-treated CNTs (Fig. 6b), in comparison with those of as-received CNTs (Fig. 6a). It is acknowledged that using concentrated acids (a mixture of H\(_2\)SO\(_4\) and HNO\(_3\) with 3:1 v:v in current case) to treat raw CNTs can introduce oxygen-containing groups to the CNT surfaces, such as carbonyl, carboxyl and hydroxyl groups or simply oxy-groups [23]. These functional groups not only improve the hydrophilicity of the CNTs, but also enhance the overall charge storage performance [24].

The pair of broad peaks on the CV of Fig. 6b indicates the oxy-groups on the surfaces of the acid-treated CNTs had undergone a reversible redox or faradaic reaction. Since the CV with the broad peaks still remains fairly rectangular, it can still be regarded as being capacitive. Thus, the acid-treated CNTs exhibited a much higher specific capacitance of 60 F g\(^{-1}\) than that of the as-received CNTs (9.1 F g\(^{-1}\)), because the charge storage process of the former
Fig. 6. CVs of (a) as-received CNT (1.3 mg) and (b) acid-treated CNT (1.5 mg) in 1.0 M KCl aqueous solution with indicated electrode capacitance averaged over the potential range \( (C_{av}) \) and specific capacitance \( (C_s) \). Potential scan rate: 10 mV s\(^{-1}\).

involved both EDL capacitance and pseudocapacitance. It is worth mentioning that Fig. 6b shows that the increased charge storage occurs mostly in the middle of the potential window. This means that in a symmetrical capacitor, the increased charge capacity would be utilised at low voltages and hence contribute little to the energy capacity. It would be more meaningful to use such acid-treated CNTs or activated carbon in an asymmetrical supercapacitor so that the charge storage can increase at higher cell voltages to give rise to larger energy capacity. In such cases, the increased charge storage capacity can be further exploited together with the unequal electrode capacitance strategy as will be discussed in the later section.

The influence of cations in the electrolyte on capacitance was also investigated by recording CVs in CaCl\(_2\) and KCl solutions, respectively. It was particular interesting to observe that the CVs of acid-treated CNTs recorded in the CaCl\(_2\) solutions generally exhibit greater currents than those in the KCl solutions (see Fig. 7a and b). These differences could be attributed to the greater affinity of the Ca\(^{2+}\) ions than K\(^+\) towards the oxy-groups on the CNT surfaces [25]. In fact, it has been mentioned in some early studies [26, 27] that the so-called Calcium Ionophore A23187 or Antibiotic A23187 has high selectivity for binding and
Fig. 7. CVs of (a,b) acid-treated CNT recorded in (a) 1.0 and (b) 0.5 M solutions of CaCl$_2$ (red curves) and KCl (black curves) [25], and of (c,d) Haycarb carbon both recorded in 0.5 M solutions of CaCl$_2$ (red curves) and KCl (black curves). Note that the electrode loading and scan rate are different and respectively indicated along with some of the capacitance values. Arrows indicate the direction of the potential scan. The CNTs were treated using the same method as those in Fig. 6, but in a recent batch [25].

transporting the Ca$^{2+}$ ion over other cations, such as Na$^+$, K$^+$ and Mg$^{2+}$ ions in biological systems. Therefore, if we could interpret this high Ca$^{2+}$ ion selectivity of A23187 as the various oxygen containing groups in the ionophore molecule formed a Ca$^{2+}$ ion binding site/cavity as illustrated in Fig. 8a, the unique interaction between Ca$^{2+}$ and the CNT surface oxy-groups can be similarly explained in the current case. As schematically illustrated in Fig. 8b, it can be assumed that the oxy-groups on CNT surfaces have high selectivity or affinity.
for Ca\(^{2+}\) ions. This assumption is in agreement with many previous studies involving interactions between CNTs and Ca\(^{2+}\) ions [26-29]. It was also reported that according to molecular dynamic simulation, the interaction between the carboxylate group on CNTs and Ca\(^{2+}\) ions are more favourable over K\(^+\) and Na\(^+\) ions [28]. As a result of the greater affinity of the surface oxy-groups with Ca\(^{2+}\) than with K\(^+\), about 30% capacitance increase was observed as derived from the CVs (red curves) in Fig. 7a and b. Moreover, in an experimental study, as-prepared CNTs (which must contain surface oxy-groups) were demonstrated to be effective scaffold materials for osteoblast proliferation and bone formation, thanks to partly the ability of the CNTs to conduct Ca\(^{2+}\) currents [29]. This means the structural arrangement of the oxy-groups on the CNT surfaces are suitable for binding Ca\(^{2+}\) via Reaction (1) below.

\[
2 \text{R-COOH} + \text{Ca}^{2+} + 2 \, e^{-} \leftrightarrow (\text{R-COO}^-)_{2}\text{Ca}^{2+} + 2 \, \text{H}_{(ad)}
\] (1)

In addition, the fact that each Ca\(^{2+}\) ion has two positive charges could be another important cause for the increase of the specific capacitance. To compare a 1:1 type electrolyte with a 1:2 type, e.g. KCl vs. CaCl\(_2\), a simple way was adopted in this work to halve the concentration of the latter, i.e. 1.0 M KCl vs. 0.5 M CaCl\(_2\). It is obvious that the CV current observed in 0.5 M CaCl\(_2\) (red curve in Fig. 7b) was greater than that observed in 1.0 M KCl (black curve in Fig. 7a), leading to a higher specific capacitance. It is noticeable that the CNT electrode was much more sensitive to replace KCl with CaCl\(_2\) than to the doubled electrolyte concentration as the concentration change (from 1.0 to 0.5 M) had a fairly small influence on the CV current and shape. This is indicative of the charge (or ion) storage in the electrode being determined mainly by the available sites on the CNT surface for ion sorption, instead of the number of ions in the electrolyte [25].

However, when similar investigations were carried out using a Haycarb carbon electrode, the greater current previously observed in CaCl\(_2\) solutions became insignificant (see the unpublished results shown in Fig. 7c and d). This is because there was much less oxy-groups
in the Haycarb carbons and the binding cavity formed between the carbon particles may not be compatible with the Ca\(^{2+}\) ion in terms of size and shape.

Fig. 8. Schematic illustration of the binding of a Ca\(^{2+}\) ion by the oxy-groups in the cavity formed (a) in the Calcium Ionophore A23187 molecule, and (b) in between the acid-treated CNTs [25].

On the other hand, for sub-zero temperature applications, organic electrolytes are commonly used [30-32], but they have lower conductivity and are more expensive than aqueous electrolytes. Mixing water with an additive may lead to low temperature benefits. The binary phase diagram of water and formamide (FA) is shown in Fig. 9a with a eutectic temperature of \(-46\) °C [25]. The binary phase diagrams of water and some simple salts are presented Fig. 9b, showing eutectic temperature below \(-50\) °C on the CaCl\(_2\)-H\(_2\)O curve. Thus, by dissolving CaCl\(_2\) into the FA-H\(_2\)O mixture, a novel organoaqueous electrolyte was prepared and proven to be able to work at temperatures below \(-60\) °C [25].

Fig. 10a and 10b compare the CVs of the CNT electrode measured in different electrolytes at 20 °C and \(-60\) °C, respectively. It can be seen that the solution of CaCl\(_2\) in FA is not a good electrolyte, but addition of FA into water does not affect very much the electrochemical behaviour. Particularly, Fig. 10b confirms the CaCl\(_2\)-FA-H\(_2\)O mixture to be
Fig. 9. (a) The formamide-water phase diagram. The dashed lines are presented as a vision guide [25]. (b) Phase diagrams of CaCl$_2$, MgCl$_2$, NaCl, CaMg$_2$Ac$_6$ (CMA), and KAc, Ac$=$CH$_2$COO$^-$ [33].

Fig. 10. (a,b) CVs and (c) $C_t/C_{20}$ (= 100×capacitance at designated temperature / capacitance at 20$^\circ$C) of CNT electrode at (a) 20$^\circ$C and (b) −60$^\circ$C in the CaCl$_2$ solution of FA (1.0 M), H$_2$O (2.0 M) and mixed FA-H$_2$O (1:1, v:v, 2.0 M). Potential scan rate: 100 mV s$^{-1}$ [25].
very good for low temperature uses. Fig. 10c presents the capacitance retention of the CNT electrode as derived from the rectangular CVs in Fig. 10a and 10b. At 20 °C, the measured specific capacitance, \( C_s \), was almost the same in CaCl\(_2\)-H\(_2\)O (108.3 F g\(^{-1}\)) and CaCl\(_2\)-FA-H\(_2\)O (107.6 F g\(^{-1}\)), whereas at -60 °C, \( C_s \) in CaCl\(_2\)-FA-H\(_2\)O was still 72.7 F g\(^{-1}\) (67% capacitance retention), which is more than double of the value measured in CaCl\(_2\)-H\(_2\)O. Hence, it can be concluded that the CaCl\(_2\)-FA-H\(_2\)O mixture is an excellent electrolyte for capacitive charge storage in partially oxidised CNTs at sub-zero-temperatures.

3. Capacitive faradaic charge storage

Charge storage in pseudocapacitive materials involves electron transfer reactions and hence the reduction or oxidation changes in the electrode materials, which are also referred to as faradaic reactions or processes. However, not all faradaic processes can contribute to pseudocapacitive behaviours. Despite the electron transfer reactions in a rechargeable battery or a fuel cell are also faradaic in nature, the respective current-voltage relations in these electrochemical devices differ significantly from that in a capacitor [9]. Such energy storage and discharge can be regarded as non-capacitive faradaic processes. The performance difference between capacitive and non-capacitive faradaic processes could be the effect of localised and delocalised valence electrons [4, 34].

3.1. Reversible non-capacitive faradaic process in electrode coatings

Non-capacitive faradaic process can result from localised reversible electron transfer reactions in a redox active coating on an electrode. When this happens, the CV should be a bell-shaped curve as displayed in Fig. 11a [3, 34]. The potential of the current peaks on the CV is determined by the redox potential of the material. The current response towards the potential change can be expressed as Eq. (10), which is derived from the Nernst and Randles-
Sevcik equations [35].

\[
i = \frac{n^2 F^2 A \Gamma \nu \exp \left[ \left( E - E^o \right) \frac{nF}{RT} \right]}{RT \left[ 1 + \exp \left[ \left( E - E^o \right) \frac{nF}{RT} \right] \right]}
\]  

where \( n \) is the number of electrons transferred between the reduced and oxidised sites, \( F \) the Faraday constant, \( A \) the area of electrode, \( \Gamma_i (=\Gamma_o + \Gamma_t) \) the total surface covered by the reduced and oxidised sites and \( \nu \) is the potential scan rate, \( R \) the ideal gas constant and \( T \) the absolute temperature. Due to the strong dependence of the electrode reaction on potential, the GCD plots in Fig. 11b show potential plateaux over a narrow potential range. Nevertheless, the faradaic reaction took place in the electrode materials is reversible as the shapes of the CVs and GCDs are symmetrical in horizontal and vertical orientations, respectively [4].

The cause for CVs or GCD plots in Fig. 11 are not rectangular or triangular like those shown in Fig. 2 may be explained from a view point of electron energy states. As electron energy states are separated in isolated molecules or insulators, the energy levels of the filled and vacant states are singular and the process that electron inject to (or remove from) an

![Graph A](image1.png)

**Fig. 11.** (a) CVs at indicated potential scan rates and (b) GCD plots at indicated constant currents of \( i_a > i_b > i_c \) for a reversible faradaic reaction with localised electron transfer to and from isolated redox sites on the electrode [4].
electrode may take place in very well separated redox active sites. In other words, because these redox centres are non-interactive or electronically isolated from each other but have equal or fairly close energy states, electron acceptance or donation occurs at potentials very close to each other, resulting in the peak shaped CV in a narrow potential range as shown in Fig. 11a. However, if these redox active sites can interact with each other due to either short separations or good electronic conductivity or both, their energy states can merge into a broad band with negligibly small differences between the neighbouring states. Such a situation corresponds to the conduction band in semiconductors, including most transition metal oxides, and is also comparable with the electron delocalisation in conjugated chemical bonds, as in electronically conducting polymers (ECPs), resulting from overlapping electron orbits between neighbouring atoms. As a result, electron transfer into (or from) each energy state in this broad band becomes continuous over a wide range of potentials, which is responsible for the constant current flow and hence the rectangular CV.

3.2. Pseudocapacitance from delocalised electrons

The above analysis is derived from the band theory for semiconductor [35], which can also be applied for qualitative description of the capacitive faradaic process for charge storage in pseudocapacitive materials. Fig. 12a illustrates the distribution energy states in three different materials, namely, metals, semiconductors, and insulators. Electron energy states of metals are completely overlapped, enabling free electron mobility in response to an electrical field. Unfortunately, no capacitance can be achieved since charges cannot be stored in a geometric location. Within an insulator, redox active sites are unable to interact with each other electrically and hence behave like individual molecules. This means localised electron transfer from the electrode substrate to each of the redox active sites in a narrow potential range and hence a current peak on the CV. Such processes are shown in Fig. 12b on the CV.
of an electro-deposited ferrocenyl polymer coating on a Pt disc electrode in an organic electrolyte [36]. The large peak couple, a2/c2, results from the ferrocenyl groups, and the smaller couple, a3/c3, from the polymer. As previously mentioned, a narrow peak would form on the CV if the non-capacitive faradaic electrode reaction takes place in which electron transfer reaction occurs at separated sites (e.g. the ferrocenyl groups) with energy states very comparable to each other.

As shown in Fig. 12a, semiconductors are featured by their slightly separated filled valence band and the vacant conduction band. When the applied potential is suitable for electron excitation, electrons in the valence band can jump over to the conduction band and become delocalised. At the same time, mobile holes are left in the valence band. Thus, although the energy gap between the broad valence and conduction bands is small, the continuous electron dislocation and hole production over the relatively wider potential range give rise to the constant capacitive current flow as shown in Fig. 12c. These CVs were recorded in 0.5 M LiCl using an electrode coated with thin films of MnO$_2$ with different thicknesses.[37] Due to electron dislocation in semiconductor MnO$_2$ coatings conducted amongst redox sites which interacted actively and continuously with each other between neighbouring states, rectangular CVs were observed. This type of charge storage process is the so called capacitive faradaic process.

It is noteworthy that mixed forms of the above two types of charge storage mechanisms are also observed in many cases with the development of advanced nanomaterials for supercapacitor applications. For instance, it has been reported that the electro-deposited polyaniline (PAN) nanofibrils were able to store charges relying on both pseudocapacitance and EDL capacitance, giving rise to a maximum specific capacitance for PAN.[38] It was also suggested that both charging processes were carried out in parallel. In other words, non-faradaic capacitive, capacitive faradaic and non-capacitive faradaic reactions have all
proceeded in PAN and made contributions to the overall charge storage.

Fig. 12. (a) Schematic illustration of the band theory [34]. (b,c) CVs of (b) an electrodeposited ferrocenyl bithiophenic polymer coating on Pt disc electrode in an acetonitrile electrolyte at 20 mV s\(^{-1}\) [36] and (c) MnO\(_2\) thin films of increasing thickness formed by redox deposition on graphite electrode in 0.5 M LiCl at 25 mV s\(^{-1}\) [37].

In the work, the PAN nanofibrils were potentiostatically deposited at 1.0 V vs. Ag/AgCl on a platinum disc electrode (1.6 mm in diameter) from a solution containing 0.25 M aniline monomer and 1.0 M HCl. The polymerised PAN can be presented by the chemical formula in Fig. 13a, showing the faradaic charge storage stoichiometry of PAN. The \(\alpha\) number is used to express the fraction of the charge of an electron that each aniline monomer unit shares, e.g. 0.5 electron per monomer (\(\alpha = 0.5\)) is needed to oxidise the polymer from its fully reduced state (leucoemeraldine) to the partially oxidised state (emeraldine). Upon oxidation, the polymer is doped by counter anons, A\(^-\), to balance the positive charges in the polymer chains.
Fig. 13. (a) Faradaic charge storage stoichiometry of PAN, (b) SEM image and (c) CVs of PAN nanofibrils in 1.0 M HCl at 20 mV s\(^{-1}\) in different potential ranges [38]. (d) \(C_s\) or the \(Q/(\Delta E \cdot W)\) ratio derived from the CVs in (c) with indicated potential ranges.

along which the electrons are delocalised amongst several neighbouring monomers. The prepared PAN nanofibrils had a thickness in the range of 50 – 150 nm and were interconnected to form a porous structure as shown in Fig. 13b.

CVs were also recorded in 1.0 M HCl in different potential ranges to characterise capacitive behaviours of the PAN nanofibrils. It can be seen in Fig. 13c that the CV shape changes significantly with positively shifting the negative potential scan limit and becomes rectangular with diminishing or disappearance of peak a1. This indicates the PAN behaved like a capacitor within a narrow potential range in which reversible conversion between leucoemeraldine and emeraldine was the major electrode reaction [38]. With further
extension of the potential window, PAN however behaved like a battery since the polymers started experiencing irreversible changes caused by the oxidation to the fully oxidised state, pernigraniline. Again, due to the non-capacitive or battery-like behaviour of PAN, the peak shaped CVs should not be used for the capacitance measurement as the large peak current will lead to a misleading ultra-high capacitance value [39].

In contrast, the specific capacitance of PAN can be calculated as below based on the rectangular CVs in Fig. 13c after combining faradaic and double layer contributions,

\[ C_s = \frac{Q}{\Delta E \cdot W} = \frac{\alpha \cdot F}{\Delta E \cdot M} + AC_{dl} \]  \hspace{1cm} (11)

where \( F \) is the Faraday constant, \( \Delta E \) the potential range, \( M \) the molecular mass of the monomer unit in the polymer, \( A \) the specific area (m\(^2\) g\(^{-1}\)) and \( C_{dl} \) the double layer capacitance per unit true area of the material (F m\(^{-2}\)). In accordance with the redox mechanism of PAN, \( \alpha \) varies with \( \Delta E \), but the relation may not be linear. Thus the maximum value of \( C_s \) can appear at smaller values of \( \alpha \) and \( \Delta E \) [38, 40], where the reversible redox reactions (capacitive faradaic reactions) of PAN predominate. In this situation, electron delocalisation occurs along the polymer chains to enable high conductivity and pseudocapacitance. However, it is worth pointing out that the \( C_s \) values derived from the rectangular CVs of PAN are generally smaller than those \( Q/(\Delta E \cdot W) \) ratios measured over the wide potential range in which battery-like behaviour dominated as shown in Fig. 13d. This is due to the insulating-to-conducting phase transition in the polymer leading to the large current of peak a1 [41].

Interestingly, measuring capacitance based on a non-capacitive faradaic charge storage process does not come singly but in pairs, giving misleading results. It was reported that combination of an activated carbon electrode with a redox electrolyte containing \( \text{CuCl}_2 \) could increase significantly the charge storage capacity of activated carbon, compared with a
conventional inert electrolyte as shown in Fig. 14a [42]. As expected, a huge current peak was observed because in the acidic media (1 M HNO₃), CuCl₂ reacted with the oxy-groups on the surface of the activated carbon electrode by self-reduction to form an additional redox layer of CuCl. Unfortunately, the CVs in Fig. 14a were misused to derive an ultrahigh but misleading specific capacitance of 4700 F g⁻¹. The related GCD plots (Fig. 14b) are akin to those of rechargeable batteries other than the triangular plot of a capacitor as shown in Fig. 2b. Similar to the aforementioned cases, although this electrode-electrolyte reaction was faradaic in nature, it took place at fixed or localised redox activated sites (C=O bond on the carbon surface) which are not interactive with each other to generate pseudocapacitance.

However, in the paper, the performance characteristics were still calculated based on the peak shaped CVs regardless the origin of the charge storage, giving rise to ultrahigh values of 4700 F g⁻¹ for specific capacitance of the electrode and 163 Wh kg⁻¹ for average energy.
capacity. These results particularly need to be considered since the charging energy had been underestimated and discharging energy been overestimated in comparison with the symmetric triangular GCD plot of an ideal capacitor (the red dashed plot in Fig. 15). This is because the calculated energy is the area under the dashed straight lines, whereas, the actual energy is the area under the curved solid lines. The charging energy is larger than the discharging one, reflecting the electrode reaction being not fully reversible and hence only a portion of the

![GCD plot](image)

**Fig. 15.** Schematic comparison of the GCD plot of a quasi-reversible faradaic reaction such as the case in Fig. 14 (black solid lines) with that of an ideal capacitive faradaic reaction (red dashed lines). The area under the potential curve of the GCD is proportional to the energy consumed (or released) during charging (or discharging) [43].

energy used during the charging period was released during the discharging period. In this case, the energy efficiency considered as the ratio of discharging energy to charging energy is far smaller than 1, which is not ideal for energy storage.

Different charge storage mechanisms have been briefly introduced, including non-faradaic EDL storage or capacitive faradaic pseudocapacitance and non-capacitive faradaic battery-like behaviours. However, classification of individual charge storage mechanisms has become rather ambiguous since the emergence of advanced composite electrode materials (e.g. the composites of carbon materials with ECPs [16] or transition metal oxides [44]) or
the application of a novel redox active electrolyte [42]). Confusions therefore appeared in the literature [39, 45-47] in terms of understanding of the origin of pseudocapacitance and determination of capacitance. Although Conway had stated an average pseudocapacitance of a battery could be determined by dividing the withdrawable charge ($\Delta Q$) with the voltage difference ($\Delta E$) in 1991, he simultaneously emphasised that the discharge characteristics of $Q$ vs. time would not be similar as those of a capacitor, which is not very useful and informative and could be misleading [48]. Hence, it is essential and crucial to carefully define and differentiate the charge storage process according to the electrochemistry of the electrode materials and from the viewpoint of the semiconductor or delocalisation band theory. This is certainly significant and helpful to precisely evaluate the performance of supercapacitors and beneficial for designing practical devices for energy storage, such as supercapattery.

4. Energy capacity enhanced by unequal capacitance

4.1. Unequal capacitance theory

It is well known that supercapacitors are high power energy storage devices, but relatively low in energy capacity compared to batteries. There is a strong desire for improving the energy capacity of supercapacitors. In theory, this improvement can be achieved by either increasing the capacitance or expending the working voltage of a supercapacitor based on Eq. (5), where $\varepsilon$ represents the energy stored in a supercapacitor, $C$ capacitance and $U$ voltage.

For symmetrical supercapacitors, capacitance equalisation of the two electrodes is necessary to maximise material utilisation and specific energy storage capacity, because the positive and negative electrodes have the same working potential range. Practically, one should also consider if the open circuit potential (OCP) of the material is at the centre of the potential range [49], because the OCP does not always match the potential of zero charge.
(PZC) with reference to Fig. 16. However, it may not be the same case in asymmetrical supercapacitors. The amount of charge, $Q$, stored in the positive and negative electrodes in an asymmetrical supercapacitor must be the same and is governed by Eq. (12) to (14),

$$Q = C_P U_P = C_N U_N \text{ or } U_N = U_P C_P / C_N$$  \hspace{1cm} (12)

$$C = \frac{C_P C_N}{C_P + C_N} = \frac{C_P}{1 + C_P / C_N}$$  \hspace{1cm} (13)

$$U = U_P + U_N = U_P \left(1 + C_P / C_N\right)$$  \hspace{1cm} (14)

where $C_P$ and $C_N$ are the capacitance, and $U_P$ and $U_N$ are the “working potential range” of the positive and negative electrodes in the cell. Note that $U_P$ or $U_N$ is defined as the difference between the more positive (or less negative) and less positive (or more negative) potential limits of the working potential range. It means that $U_P$ or $U_N$ is always positive, disregarding the electrode’s polarity.

Another upper and lower limit, “exploitable potential range” for positive or negative electrode, $U^o_P$ or $U^o_N$, is set by any irreversible electrode reactions in a supercapacitor cell, such as solvent decomposition and over-oxidation or over-reduction of the electrode materials. Thus, $U_P \leq U^o_P$ and $U_N \leq U^o_N$. Eq. (15) describes the relationship among $U$, $U_P$, $U_N$, $U^o_P$ and $U^o_N$. (Note: $U^o_P = \text{CPR}(\text{+})$ and $U^o_N = \text{CPR}(\text{–})$ in Fig. 16.)

$$U = U_N + U_P \leq U^o_P + U^o_N$$  \hspace{1cm} (15)

Along the potential axis, $U_P$ and $U_N$ are next to each other, but $U^o_P$ and $U^o_N$ could be either connected or overlapping partially or completely. From Eq. (12), it can be found that if $C_P$ increases, $U_P$ will reduce and the same correlation applies to $C_N$ and $U_N$. For a symmetrical supercapacitor, the above situation is avoided because of the equalised electrode capacitance (i.e. $C_N = C_P$). In the case of an asymmetrical supercapacitor, the two electrodes have likely different values for $U^o_P$ and $U^o_N$. When the positive and negative electrodes are equalised in capacitance, the cell voltage will be limited either by the positive electrode if $U^o_P < U^o_N$, or
by the negative electrode if $U'_p > U'_N$. In cases where $U'_p$ and $U'_N$ overlap, they should be replaced by $U_p$ and $U_N$ respectively. However, the voltage of an asymmetrical supercapacitor can also be extended by increasing the capacitance of the “cell voltage limiting electrode”, leading to the increase of the energy stored by the supercapacitor as indicated in Eq. (16) which is derived from Eq. (5) and (12) to (14).

$$E = \frac{1}{2} C U^2 = \frac{1}{2} C_p U_p^2 (1 + C_p/C_N)$$

(16)

**Fig. 16.** A model of supercapacitor voltage. Schematic illustration of supercapacitor maximum charging voltage (MCV), potential of zero voltage (PZV) and electrode capacitive potential range (CPR) [50]. CPR(+) = $U'_p$ and CPR(−) = $U'_N$ in Eq. (15).

It has been proven that the positive to negative electrode capacitance ratio ($C_p/C_N$) plays an important role in the improvement of a practical asymmetrical supercapacitors without changing the materials [51]. Along with $C_p/C_N$, the capacitive potential range (CPR) and potential of zero voltage (PZV) may affect the maximum charging voltage (MCV) and stored energy of the supercapacitor. **Fig. 16** shows a model of supercapacitor voltage [50]. The lower and upper potential limits of the negative and positive electrodes are denoted as $E_{N1}$,
$E_{N2}$, $E_{P1}$ and $E_{P2}$, respectively. The CPR of the negative or positive electrode is $E_{N2} - E_{N1}$ or $E_{P2} - E_{P1}$, respectively. PZV is defined as the equal electrode potential when a supercapacitor is fully discharged and its voltage is zero if $E_{N2} \geq E_{P1}$. PZV affects the coulombic efficiency and the MCV of a supercapacitor, but is difficult to predict in practice. As shown in Fig. 16, if the positive electrode is the “cell voltage limiting electrode” as the MCV of the supercapacitor is determined by $E_{P2}$, then $PZV - E_{N1} > E_{P2} - PZV$, and the capacitance of the two electrodes are equal, the MCV would be $(E_{P2} - PZV) \times 2$.

4.2. Asymmetrical supercapacitor

A few asymmetrical supercapacitor examples have proved the practicality of the unequal capacitance theory [50, 51]. An asymmetrical supercapacitor composing of a PAN-CNT nanocomposites positive electrode and a CMPB (Cabot Monarch 1300 pigment black) negative electrode has been investigated [51]. PAN has very high theoretical specific capacitance (750 F g$^{-1}$) [40] and excellent charging-discharging cycle stability [16], but it has a small exploitable potential range of 0.7 - 0.8 V in acidic aqueous electrolytes [16]. If PAN is used in a symmetrical supercapacitor, its energy capacity and power capability are compromised, as indicated in Eq. (5), despite of the polymer’s ultrahigh specific capacitance. In the previous research [51], several PAN-CNT positive electrodes were prepared at deposition charges of 166 to 312 mC, corresponding to $C_P/C_N$ ratios of 0.8 to 1.5. Fig. 17 shows the CVs of these PAN-CNT (+)|HCl (1.0 M)|CMPB (−) asymmetrical supercapacitors. In this supercapacitor, the positive electrode PAN-CNT (+) is the “cell voltage limiting electrode” [51]. In Fig. 17, the charging peaks can be identified on all these CVs near the high voltage ends and the peak voltage shifted to higher values with increasing $C_P/C_N$ ratio. When the $C_P/C_N$ ratio changes from 1.0 to 1.5, the cell capacitance only decreases by 7%, while the specific energy of the cell increases over 80%. These improvements are remarkable,
as there was only a small increase of 12.9% in the total mass of the electrode materials.

Fig. 17. CVs of asymmetrical supercapacitors with a 0.3 mg CMPB negative electrode and a PAN-CNT positive electrode at indicated $C_P/C_N$ ratios without the charging current peek. Scan rate: 20 mV s$^{-1}$ [51].

Similar asymmetrical supercapacitor composing of polypyrrole-CNT (PPy-CNT) composites positive electrode and CMPB negative electrode have been fabricated and tested. The relative research is still ongoing and unpublished. The preliminary data indicates that the asymmetrical cell can improve the cell voltage form 1.0 to 1.4 V and cell specific energy from 13.0 to 16.6 J g$^{-1}$ comparing to the symmetrical cell. Recently, graphenes are popularly used to form composite with conducting polymers that could increase supercapacitor cell voltage when combined with suitable counter electrode and electrolyte [52].

4.3. Unit cell supercapacitor

Conventionally, a symmetrical aqueous carbon-carbon supercapacitor is perceived to show a narrower working voltage than an asymmetrical supercapacitor. However, recent study suggested that a unit cell supercapacitor consisting of CMPB can reach beyond 1.9 V in
cell voltage, and retain over 85% of the initial capacitance through 10,000 cycles with the cell becoming stable after the 1,000th cycles [53]. Fig. 18 shows the CVs of CMPB in 0.3 M K$_2$SO$_4$ solution at a scan rate of 5 mV s$^{-1}$ [53]. Basically, the negative current reflects the rate of absorption/adsorption of hydrated cations, H$_3$O$^+$ and K$^+$, and desorption of anion, SO$_4^{2-}$, during the negative polarisation. The positive current corresponds to the rate of desorption of the mentioned cations and absorption/adsorption anions in the course of positive polarisation.

It is also noticed that there is small positive current humps between 0.4 and 0.5 V on the CVs when the negative potential limit reached beyond −1.0 V, where a fast increasing negative current appeared due to the formation of hydrogen on the electrode. Actually, the current humps locate in a very positive potential range. It suggests that the hydrogen is trapped in the nanopores of CMPB at high negative polarisation.

The SEM and TEM images of the CMPB used in the research are shown in Fig. 19 and reveal micrometer sized aggregates of uniform onion-like spherical nanoparticles (10 nm) [54]. This nanoporous structure could also make contribution to the extension of the negative potential limit away from the theoretical value of water decomposition [55].

Fig. 18. CVs of the conventional three-electrode cell using the CMPB-PTFE (2.5 mg) coated graphite disc working electrode in 0.3 M K$_2$SO$_4$ solution over different working potential windows. Scan rate: 5 mV s$^{-1}$ [53].
Fig. 19. TEM and SEM (insert) images of the CMPB [54].

Fig. 20 shows the CVs of CMPB-PTFE (1.0 mg) in the conventional three-electrode cell with the PZV of 50 mV and potential ranges with unequal electrode capacitance [53]. The $C_p/C_N$ ratio comes to 4/3 and has successfully extended the negative electrode potential windows towards a more negative range. Meanwhile, the positive electrode potential windows kept the same, but the cell voltage increased from 1.6 V ($C_p/C_N = 1$) to 1.9 V ($C_p/C_N = 4/3$). The energy capacity and capacitance have been compared between equal and unequal electrode capacitance CMPB supercapacitors. Fig. 21 shows the capacitance retention of supercapacitor with different $C_p/C_N$ ratios [53]. It can be found that the supercapacitor with the $C_p/C_N$ ratio of 4/3 performs better than the one with the ratio of 1. The calculation of the cell capacitance and specific energy capacity indicate that there will be only 2% decrease of the cell capacitance but 38% increase of the cell specific energy capacity. This is a typical zero cost engineering example.
Fig. 20. CVs of CMPB-PTFE (1.0 mg) in the conventional three-electrode cell with the PZV of 50 mV and potential ranges with unequal electrode capacitance [53].

Fig. 21. Capacitance retention (C/C\textsuperscript{0}) as a function of the number of charge-discharge cycles of sandwich-type supercapacitors with the cell voltage limit indicated [49].

5. Supercapattery

Another widely applied approach to increasing the energy capacity of supercapacitors is using hybrid configurations that take the advantages of both the supercapacitor and battery. The particular design involves combination of a supercapacitor electrode with a battery electrode, such as the so called lithium ion capacitor [56] and the more general term supercapattery (= supercapacitor + battery) [50, 57].

Theoretically, the hypothetical battery of lithium metal and fluorine gas (Li-F battery)
would output a cell voltage about 6.1 V and offer a specific energy content of 6304 Wh/kg at room temperature. This value sets the upper limit of specific energy for all batteries. For Li-ion batteries (e.g. Li$_x$C$_6$ | Li$_{1-x}$CoO$_2$), the theoretical specific energy is 552 Wh/kg at 3.5 V. If one considers a supercapattery with a lithium metal negative electrode and a supercapacitor positive electrode of 400 F/g, the specific energy value would be 625 Wh/kg for the cell voltage to vary from 3.5 V to 1.0 V. In this calculation, the mass of the lithium metal is negligible because the specific charge capacity of lithium is much larger than that of the supercapacitor electrode. Thus, in theory, supercapattery can possess higher specific energy than both battery and supercapacitor, and should be able to supply this specific energy at a power output almost as high as supercapacitor with suitable electrode and cell designs.

Nanostructured carbons, such as activated carbon, CNT and graphene represent the primary material choices for the polarisable capacitor-like electrode (i.e. capacitive non-faradaic) owing to their large surface area, porosity, stability over a wide potential window, as well as intrinsically low electrical resistance. Pseudo-capacitive materials including MnO$_2$, RuO$_2$ and conducting polymers can also be used to gain higher electrode capacitance which however also means narrower potential windows. As to the battery-like electrode, a much wider choice of materials exists, from lead acid battery to metal/air systems, but metal compounds such as SnO$_2$, MnO$_2$ and LiFePO$_4$ are more practically and commercially available. It should be mentioned that combining the supercapacitor and battery electrodes into one device could result in performance very close to that of the supercapacitor [1, 50]. This is demonstrated by the calculated GCD plots shown in Fig. 22 for a cell with a lithium metal negative electrode and an activated carbon positive electrode. In the calculation, the potential range of the activated carbon is set between 0.5 V and 4.5 V vs. Li/Li$^+$ to avoid interaction of the carbon with lithium metal. With overall cell performance being comparable to that of a supercapacitor, the analysis and comparison can follow those capacitor equations.
Examples of such capacitive cell behaviour are also reported in the literature [58, 59]. On the other hand, the combination of battery and supercapacitor electrodes may also present battery like features as that shown in Figs. 14 and 15. In such cases, data presentation and analysis should follow those for batteries. Particularly, specific charge capacity (e.g. mAh g$^{-1}$) should be reported, instead of specific capacitance (F g$^{-1}$). For hybrid devices showing rechargeable battery behaviour, the term supercabattery is recommended [50].

![Graph](image)

**Fig. 22.** Calculated potentials of the positive activated carbon electrode (blue solid line) and the lithium metal negative electrode (solid black line) and the cell voltage (dashed red line) for a “Li | activated carbon” supercapattery.

Recent studies reveal that an aqueous supercapattery consisting of a MnO$_2$ positive electrode and a Li/LISICON/PEO-LiTFSI/Li$^+$ negative electrode had achieved a specific energy value of 114 Wh kg$^{-1}$ with a 4.3 V cell voltage [60]. The specific energy can come to 520 Wh kg$^{-1}$ with a 3.8 V cell voltage, if the positive electrode is replaced by RuO$_2$ [60]. However, the current density of the supercapattery discussed above is 0.255 mA cm$^{-2}$, which is still not ideal for high power application. Our recent supercapattery work based on an activated carbon (110 F g$^{-1}$) positive electrode and a Li/Li$^+$ negative electrode using an ionic...
liquid electrolyte can achieve a specific energy value of 200 Wh kg\(^{-1}\) at a 1.0 mA cm\(^{-2}\) current density. The relative work is still ongoing and will be published elsewhere [61].

6. Conclusions

In this paper, we have proposed to define and differentiate capacitive and non-capacitive faradaic processes for charge storage in supercapacitors according to the band theory which the origin of pseudocapacitance has been correlated with electron transfer to or from the conduction band of semiconductor type materials. The principle and performance governing equations of conventional capacitors have been presented and applied to differentiate between batteries and supercapacitors based on the shapes of CVs and GCDs. It is confirmed that pseudocapacitance is featured by rectangular CVs and triangular GCDs and peak-shaped CVs and non-linear GCDs are features of battery behaviour which should not be used for measurement of capacitance values and energy estimation.

We have also found a unique affinity between Ca\(^{2+}\) ions and the oxy-groups on the surfaces of CNTs, leading to a significant capacitance improvement. The specific capacitance of the CNTs measured in CaCl\(_2\) based electrolytes was generally higher than that in more commonly used alkali salts based electrolytes. This understanding is new to supercapacitor research and could have a significant impact on development of supercapacitors based on CNTs, graphenes and/or activated carbons which are all enriched with various surface oxy-groups. In particular, after mixing with FA and H\(_2\)O, the resultant organoaqueous CaCl\(_2\) electrolyte enabled the CNTs to store charges at both room and sub-zero temperatures below −60 °C without a significant capacitance loss.

On the other hand, energy capacity enhancement by unequal capacitances has been demonstrated in theory and examples. It clear shows that the phenomena of the potential of zero voltage (PZV), capacitive potential range (CPR) and positive to negative electrode
capacitance ratio ($C_P/C_N$) affect the performance of supercapacitors. It also shows that for energy enhancement, the capacitance unequalisation approach could be very low or almost zero in engineering cost. The recent development in supercapattery that utilises both capacitive and non-capacitive charge storage mechanisms (i.e. hybrid devices combining supercapacitor and battery electrodes) reveals much increased energy capacity comparable or higher than those existing rechargeable batteries, including the lithium ion battery. These initial studies have demonstrated a very promising prospect of supercapattery.

Acknowledgement

This work received funding from the EPSRC (GR/R68078/02), E.ON AG (International Research Initiative --- Energy Storage 2007), and the Ningbo Municipal Government (3315 Plan and the IAMET Special Fund, 2014A35001-1). Responsibility for the content of this paper lies with the authors.

Reference


[27] H. Wang, H. Nie, Y. Wang, Y. Chang, E. Dong, S. Yang, Purifying carbon nanotube polluted water involves adding calcium chloride aqueous solution with carbon nanotube aqueous solution, mixing mixture and storing it to sediment particles, and obtaining colorless transparent water, Univ Shanghai.


based electrolyte in supercapacitors. A safe and wide temperature range electrolyte,


